ABSTRACT

BUCKLING BEHAVIOUR OF SINGLE-WALLED CARBON NANOTUBES UNDER AXIAL LOADING

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We investigate the behaviour of a single-walled Carbon Nanotube under axial compressive line load applied at both edges. The expected buckling response is studied by application of a molecular computation model. We formulate a global potential and search for its minimum to obtain the equilibrium configuration. The critical nanotube diameter, when local shell buckling occurs, is measured with two parameters: the

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value of compression loading and tube diameter.

INTRODUCTION

Carbon nanotubes (CNTs), after their discovery in 1991 by Iijima [6], attracted a lot interest from researchers working mainly in the new field of nanotechnology. Among unusual properties of these new structures there are mechanical properties [3, 8]. For instance the Young's modulus of a carbon nanotube is expected to be several times that of the diamond (\sim 1 TPa) reaching the extreme values compared to any other material. CNTs show various electrical properties as metallic conductors or semiconductors, depending on their physical properties [12].

Their unique physical properties have lead to various applications in different areas, such as scanning of chemicals content [7], nanoelectronics [5, 9] and biosensors [13, 2], nanoelectromechanical storage systems [4]. Nanotubes were applied as nanosensors and nanoactuators. Namely, vertically aligned single-walled carbon nanotubes were used in emission and vacuum micro-electronic devices, and as nano-, bio-sensors [13] and nanoactuators [5].

In this note, motivated by potential applications of these interesting materials [11], we examine a critical value of axial loading which leads to local buckling instability for a single walled carbon nanotube.

MODEL ASSUMPTIONS

In the first step we define interactions between carbon atoms forming a tube. They can be defined by a simplified inter-atomic phenomenological potential V of Brenner-Teresoff type [1, 10]:

$$V = \sum_{i,j(>i)} V_R(r_{ij} - \overline{B}_{ij}V_A(r_{ij})$$
(1)

where *i* and *j* indicate the atomic sites, while r_{ij} is the distance between *i*-th and *j*-th carbon atoms while $\overline{B}_{ij} = (B_{ij}+B_{ji})/2$. The related nearest neighbours configuration and the applied axial forces are presented in Figure 1. The system parameters are included in Table 1.

Figures 2a and 2b show atomic bond configuration under loading for armchair carbon nanotubes of two different diameters. We show the deformation of a short piece of a carbon nanotube under uniaxial compression preserving the cylindrical symmetry. Such deformation is connected with small changes in nanotube geometry in both diameter and length.



Fig. 1. Definition of nearest neighbour bonds angle ϕ_{ijk} (a), schematic plot of uniaxial compression |F'| = |F| = F (b)

Table 1. Th	e table	show	а	set	of	parameters	used	in
calculations	[1, 10]							

Parameter	Value			
R _{ii} ^(e)	1.39 Å			
D _{ii} ^(e)	6.0 eV			
β _{ii}	2.1 Å ⁻¹			
S	1.22			
δ	0.5			
α	0.0			
R _{ii} ⁽¹⁾	1.7 Å			
R _{ii} ⁽²⁾	2.0 Å			
α	0.000200813			
C ₀	330			
d	3.5			

STABILITY ANALYSIS

In the next step we analyse the atomic potential and stability of the deformed structures (Fig. 2). The repulsive and attractive parts of the potential: V_R , V_A are defined as:

$$V_{R}(r_{ij}) = \frac{f_{ij}(r_{ij})D_{ij}^{(e)}}{S_{ij}-1}\exp(-\sqrt{2S_{ij}}\beta_{ij}\Delta r_{ij})$$
(2)

$$V_{R}(r_{ij}) = \frac{f_{ij}(r_{ij})D_{ij}^{(e)}S_{ij}}{S_{ij}-1}\exp(-\sqrt{2S_{ij}}\beta_{ij}\Delta r_{ij})$$
(3)

where:

$$f_{ij} = \begin{cases} 1, & r < R_{ij}^{(1)} \\ \left[1 + \cos\left(\frac{\pi \left(r - R_{ij}^{(1)}\right)}{\left(R_{ij}^{(2)} - R_{ij}^{(1)}\right)}\right) \right] / 2, & R_{ij}^{(1)} < r < R_{ij}^{(2)} \\ 0, & r > R_{ij}^{(2)} \end{cases}$$
(4)

$$B_{ij} = \left[1 + \sum_{k(\neq i,j)} G_i(\phi_{ijk}) f_{ijk}(r_{ij}) \exp(\alpha_{ijk} \left[\Delta r_{ij} - \Delta r_{ik}\right]\right]^{-\delta_{ij}}$$
(5)

$$G_{i}(\phi_{ijk}) = a_{0} \left\{ 1 + c_{0}^{2} / d_{0}^{2} - c_{0}^{2} / \left[d_{0}^{2} + \left(1 + \cos \phi_{ijk} \right)^{2} \right] \right\}$$
(6)

$$\Delta r_{ij} = r_{ij} - R_{ij}^{(e)} \tag{7}$$



Fig. 2. Projections of 2×2 (a) and 5×5 (b) single walled carbon nanotubes structures (armchair) under an axially directed compression on their edges with a F = 3 eV/Å. Here, red lines correspond to loaded nanotube structures while black to unloaded ones (distances are expressed in carbon lattice spaces of the unloaded tube structure)

Symbol ϕ_{ijk} defines the angle between the lines, connecting *i*-th with *j*-th atom and *i*-th with *k*-th atom. $R_{ij}^{(e)}$ is the distance between atoms at equilibrium. All the potential parameters appearing in Equations 1-7 are given in Table 1, and were adjusted from Ref. [1].

To determine the stability of the given equilibrium structure under the loading we analyze the Jacobian matrix of the total potential $V_{TOT}(r_1, r_2, ..., r_N)$ composed of the inter-atomic potential V (Eq. 1) and the work of the external axial forces (Fig. 1b) **F**=-**F**'.

To analyze the stability of atom configuration we used the condition that for a stable configuration of atoms eigenvalues of the Jacobian matrix calculated from V_{TOT} with the following matrix elements:

$$\frac{\partial^2 V_{TOT}}{\partial \xi_i \partial \eta_j} \tag{8}$$

where ξ_i , $\eta_j = x_1$, y_1 , z_1 , x_2 , y_2 , z_2 ,..., x_n , y_n , z_n

are Cartesian coordinates of all N atoms defined in the studied carbon nanotube. For a stable configuration all eigenvalues of the Jacobian matrix



Fig. 3. Schematic plot of a single atom displacement (a) and critical values of the axial force versus the unloaded nanotube diameter $D_n = 3n/\pi_-$ (for the armchair structure of the size $n \times n$) where the unit value is distance of carbon nearest atoms $R_{ij}^{(e)}$. Note that the units of critical force F_c is 1 eV/ °A= 1.602176462 nN

should be positive. In the present approximation we restricted the analysis to a single carbon atom displacement (Fig. 3a) where all surrounding atoms positions were fixed. The results of the critical force F_c are shown in Figure 3b. Note that the critical force F_c , above which the system is unstable, is growing with the size of the unloaded nanotube diameter D_n . This represents a general trend expected to occur in any cylinder structure, which is usually more stable on buckling for larger D_n .

CONCLUSIONS

We investigated stability of a single-walled (armchair) Carbon Nanotube under axial compressive line loading applied at both of its edges. By analyzing the Jacobian matrix eigenvalues we successfully estimated the critical force leading to local bifurcations. To obtain more appropriate results of F_c of a local buckling instability, one should go beyond a single atom displacement. The results of such an investigation are in preparation and they will reported in a separate paper.

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